# Exercise 11: Rotation and invariants Ian. 27-31

## Question 1

We want to demonstrate for the two-dimensional case that Hooke's law with isotropic elastic constants is indeed isotropic. Consider a 2D stress tensor  $\sigma$  and the corresponding strain  $\varepsilon$ ,

$$\sigma = \begin{bmatrix} \sigma_{xx} & \sigma_{xy} \\ \sigma_{xy} & \sigma_{yy} \end{bmatrix}, \quad \varepsilon = \begin{bmatrix} \varepsilon_{xx} & \varepsilon_{xy} \\ \varepsilon_{xy} & \varepsilon_{yy} \end{bmatrix}. \tag{1}$$

Next, consider the matrix for rotation by an arbitrary angle  $\alpha$ 

$$R = \begin{bmatrix} \cos(\alpha) & \sin(\alpha) \\ -\sin(\alpha) & \cos(\alpha) \end{bmatrix}. \tag{2}$$

The most straightforward way to demonstrate isotropy would be to rotate the elastic stiffness tensor. However, this is a fourth-order tensor and rotating it is cumbersome. Here, we take a different approach. In order to demonstrate isotropy

- 1. express  $\sigma$  in terms of the components of  $\varepsilon$ ,
- 2. rotate  $\sigma$  to find the representation  $\sigma'$  of this state of stress in the new coordinate system,
- 3. replace the components of  $\varepsilon$  in  $\sigma'$  by the components of the strain tensor  $\varepsilon'$  in the rotated coordinate system.

You should see that the constants of proportionality between stress and strain - the elastic constants - are the same in the new and the old coordinate system!

## **Solution:**

1.)

$$\sigma_{ij} = \lambda \delta_{ij} \varepsilon_{kk} + 2\mu \varepsilon_{ij} \tag{3}$$

$$\underline{\sigma} = \begin{bmatrix} 2\mu\varepsilon_{xx} + \lambda\left(\varepsilon_{xx} + \varepsilon_{yy}\right) & 2\mu\varepsilon_{xy} \\ 2\mu\varepsilon_{xy} & 2\mu\varepsilon_{yy} + \lambda\left(\varepsilon_{xx} + \varepsilon_{yy}\right) \end{bmatrix}$$
(4)

2.) Assuming that R is the matrix which, given the representation of a vector in the original coordinate system, yields the representation in the new coordinate system, we need to perform the following operation to find  $\sigma'$ :

$$\sigma' = R\sigma R^T$$
 (matrix notation), or, equivalently,  $\sigma'_{mn} = R_{mi}R_{nj}\sigma_{ij}$  (index notation). (5)

The result is

$$\sigma' = \begin{bmatrix} \cos(\alpha)^{2} \sigma_{xx} + \sin(2\alpha) \sigma_{xy} + \sin(\alpha)^{2} \sigma_{yy} & \cos(2\alpha) \sigma_{xy} - \cos(\alpha) \sin(\alpha) (\sigma_{xx} - \sigma_{yy}) \\ \cos(2\alpha) \sigma_{xy} - \cos(\alpha) \sin(\alpha) (\sigma_{xx} - \sigma_{yy}) & \sin(\alpha)^{2} \sigma_{xx} - 2 \sin(\alpha) \cos(\alpha) \sigma_{xy} + \cos(\alpha)^{2} \sigma_{yy} \end{bmatrix}$$

$$= \begin{bmatrix} \sigma'_{xx} & \sigma'_{xy} \\ \sigma'_{xy} & \sigma'_{yy} \end{bmatrix}, \text{ with }$$

$$\sigma'_{xx} = (\varepsilon_{xx} + \varepsilon_{yy}) (\mu + \lambda) + (\varepsilon_{xx} - \varepsilon_{yy}) \mu \cos(2\alpha) + 2\varepsilon_{xy} \mu \sin(2\alpha),$$

$$\sigma'_{yy} = (\varepsilon_{xx} + \varepsilon_{yy}) (\mu + \lambda) - (\varepsilon_{xx} - \varepsilon_{yy}) \mu \cos(2\alpha) - 2\varepsilon_{xy} \mu \sin(2\alpha),$$

$$\sigma'_{xy} = \mu (2\varepsilon_{xy} \cos(2\alpha) - (\varepsilon_{xx} - \varepsilon_{yy}) \sin(2\alpha)).$$
(6)

3.) To get the components of  $\varepsilon$  in terms of the components of  $\varepsilon'$ , we need to consider the reverse sense of rotation, i.e.  $\varepsilon = R^T \varepsilon' R$ . The transformation rules are the same for stress and strain, therefore the result can be obtained immediately by replacing  $\alpha \to -\alpha$ ,  $\sigma_{xx} \to \varepsilon'_{xx}$ ,  $\sigma_{xy} \to \varepsilon'_{xy}$ , and  $\sigma_{yy} \to \varepsilon'_{yy}$  in the matrix above,

$$\varepsilon = \begin{bmatrix} \cos(\alpha)^2 \varepsilon'_{xx} - \sin(2\alpha)\varepsilon'_{xy} + \sin(\alpha)^2 \varepsilon'_{yy} & \cos(2\alpha)\varepsilon'_{xy} + \cos(\alpha)\sin(\alpha)\left(\varepsilon'_{xx} - \varepsilon'_{yy}\right) \\ \cos(2\alpha)\varepsilon'_{xy} + \cos(\alpha)\sin(\alpha)\left(\varepsilon'_{xx} - \varepsilon'_{yy}\right) & \sin(\alpha)^2 \varepsilon'_{xx} + 2\sin(\alpha)\cos(\alpha)\varepsilon'_{xy} + \cos(\alpha)^2 \varepsilon'_{yy} \end{bmatrix}.$$
(7)

Inserting the components of  $\varepsilon$  in the equations for  $\sigma'_{xx}$ ,  $\sigma'_{xy}$ , and  $\sigma'_{yy}$ , we obtain

$$\sigma' = \begin{bmatrix} 2\mu\varepsilon'_{xx} + \lambda\left(\varepsilon'_{xx} + \varepsilon'_{yy}\right) & 2\mu\varepsilon'_{xy} \\ 2\mu\varepsilon'_{xy} & 2\mu\varepsilon'_{yy} + \lambda\left(\varepsilon'_{xx} + \varepsilon'_{yy}\right) \end{bmatrix}. \tag{8}$$

We can see that the elastic constants are the same in the two coordinate systems.

## Question 2 .....

Reference: Rösler, Harders, Bäker, Mechanisches Verhalten der Werkstoffe, 2nd ed, Teubner, p. 412

A component made from a polycrystalline Aluminum alloy has the yield strength of  $200~\mathrm{MPa}$  and is subjected to plane stress with

$$\sigma_{xx} = \sigma_{yy} = 155 \text{ MPa}, \quad \tau_{xy} = 55 \text{ MPa}. \tag{9}$$

- (a) Write the deviatoric stress!
- (b) Calculate the principal stresses!
- (c) Evaluate both Tresca's and von Mises' criterion to determine whether the material will yield!

### **Solution:**

(a)

$$\sigma' = \sigma - \frac{1}{3} \operatorname{Tr} \sigma \mathbf{1} = \begin{bmatrix} 155 \text{ MPa} & 55 \text{ MPa} & 0 \\ 55 \text{ MPa} & 155 \text{ MPa} & 0 \\ 0 & 0 & 0 \end{bmatrix} - 310/3 \text{ MPa} \mathbf{1} = \begin{bmatrix} 155/3 \text{ MPa} & 55 \text{ MPa} & 0 \\ 55 \text{ MPa} & 155/3 \text{ MPa} & 0 \\ 0 & 0 & -310/3 \text{ MPa} \end{bmatrix}$$
(10)

(b)

$$\sigma_1 = 210 \text{ MPa}, \ \sigma_2 = 100 \text{ MPa}, \ \sigma_3 = 0 \text{ MPa}$$
 (11)

(c) The two criteria contradict in this case,

$$\sigma_{\text{Tresca}} = \sigma_I - \sigma_{III} = 210 \text{ MPa} > 200 \text{ MPa} \implies \text{ the material will yield,}$$
 (12)

$$\sigma_{\text{Mises}} = \sqrt{\frac{1}{2} \left[ (\sigma_I - \sigma_{II})^2 + (\sigma_{II} - \sigma_{III})^2 + (\sigma_I - \sigma_{III})^2 \right]} = 181.93 \text{ MPa} \implies \text{the material will not yield.}$$
(13)

Both criteria are empirical, hence no decision can be made.

## Question 3

Reference: Cleland Foundations of Nanomechanics, Springer, p. 174.

(a) A solid is subjected to the stress given below. Find the three stress invariants and the three principal values of stress. Solve for the directions of the three principal axes.

$$\boldsymbol{\sigma} = \begin{bmatrix} 1 & 1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & 3 \end{bmatrix} \text{ Pa} \tag{14}$$

(b) A solid is subjected to the stress given below. Find the expression for the stress tensor if the coordinate axes are rotated by  $60^{\circ}$  counterclockwise about the z-axis.

$$\sigma = \begin{bmatrix} 2 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 4 \end{bmatrix}$$
 Pa (15)

(c) A solid is stressed according to the tensor below. Show that for this form, the stress tensor is invariant under rotations about the *z*-axis.

$$\boldsymbol{\sigma} = \begin{bmatrix} a & 0 & 0 \\ 0 & a & 0 \\ 0 & 0 & b \end{bmatrix} \tag{16}$$

#### **Solution:**

(a) The three invariants are

$$I_{1} = tr(\sigma) = \sigma_{xx} + \sigma_{yy} + \sigma_{zz} = 5 \text{ Pa},$$

$$I_{2} = \sigma_{xx}\sigma_{yy} + \sigma_{yy}\sigma_{zz} + \sigma_{xx}\sigma_{zz} - \sigma_{xy}^{2} - \sigma_{yz}^{2} - \sigma_{xz}^{2} = 6 \text{ Pa},$$

$$I_{3} = \det(\sigma) = 0 \text{ Pa}.$$
(17)

The principal stresses are the eigenvalues  $\lambda_i$  ( $i=1\ldots 3$ ) of  $\sigma$ , i.e. the roots of the characteristic polynomial. The invariants reappear here as the coefficients of the polynomial,

$$\lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 = 0. ag{18}$$

The roots are  $\lambda_1 = 3$  Pa,  $\lambda_2 = 2$  Pa,  $\lambda_3 = 0$  Pa. The eigenvectors  $e_i$  ( $i = 1 \dots 3$ ) are obtained by solving for the components of  $e_i$  in the equation

$$\sigma \mathbf{e}_i = \lambda_i \mathbf{e}_i. \tag{19}$$

For each i, we have three equations and three unknowns (the components of  $e_i$ ),

$$\lambda_{1} : \begin{bmatrix} 1 & 1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & 3 \end{bmatrix} \begin{bmatrix} e_{1x} \\ e_{1y} \\ e_{1z} \end{bmatrix} = 3 \begin{bmatrix} e_{1x} \\ e_{1y} \\ e_{1z} \end{bmatrix} \rightarrow \mathbf{e}_{1} = \begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}, 
\lambda_{2} : \begin{bmatrix} 1 & 1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & 3 \end{bmatrix} \begin{bmatrix} e_{2x} \\ e_{2y} \\ e_{2z} \end{bmatrix} = 2 \begin{bmatrix} e_{2x} \\ e_{2y} \\ e_{2z} \end{bmatrix} \rightarrow \mathbf{e}_{2} = \begin{bmatrix} 1 \\ 1 \\ 0 \end{bmatrix}, 
\lambda_{3} : \begin{bmatrix} 1 & 1 & 0 \\ 1 & 1 & 0 \\ 0 & 0 & 3 \end{bmatrix} \begin{bmatrix} e_{3x} \\ e_{3y} \\ e_{3z} \end{bmatrix} = 0 \begin{bmatrix} e_{3x} \\ e_{3y} \\ e_{3z} \end{bmatrix} \rightarrow \mathbf{e}_{3} = \begin{bmatrix} -1 \\ 1 \\ 0 \end{bmatrix}.$$
(20)

(b) The rotation matrix that transforms from rotated to original frame is ( $\theta=60^\circ)$ 

$$\mathbf{R} = \begin{bmatrix} \cos(\theta) & -\sin(\theta) & 0\\ \sin(\theta) & \cos(\theta) & 0\\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} \frac{1}{2} & -\frac{\sqrt{3}}{2} & 0\\ \frac{\sqrt{3}}{2} & \frac{1}{2} & 0\\ 0 & 0 & 1 \end{bmatrix}.$$
 (21)

If we right-multiply this matrix with a column vector whose components are given in the rotated frame, then we get the representation of this vector in the original coordinate system. The rotated stress tensor is computed as

$$\sigma' = \mathbf{R}^T \sigma \mathbf{R}$$
 (matrix notation),  
 $\leftrightarrow \sigma'_{ij} = R_{ki} R_{lj} \sigma_{kl}$  (index notation). (22)

Carrying out the multiplication, we find

$$\boldsymbol{\sigma}' = \frac{1}{2} \begin{bmatrix} (1+\sqrt{3}) & -(1+\sqrt{3}) & 0\\ -(1+\sqrt{3}) & (3-\sqrt{3}) & 0\\ 0 & 0 & 8 \end{bmatrix}.$$
 (23)

(c) The rotation matrix has the form

$$\mathbf{R} = \begin{bmatrix} \cos(\theta) & -\sin(\theta) & 0\\ \sin(\theta) & \cos(\theta) & 0\\ 0 & 0 & 1 \end{bmatrix}. \tag{24}$$

Rotation of the stress tensor gives

$$\boldsymbol{\sigma}' = \mathbf{R}^T \boldsymbol{\sigma} \mathbf{R} = \begin{bmatrix} a & 0 & 0 \\ 0 & a & 0 \\ 0 & 0 & b \end{bmatrix} = \boldsymbol{\sigma}, \quad \text{independent of } \theta.$$
 (25)

Q.E.D.